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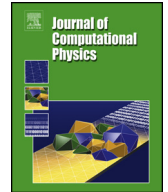
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Finite difference methods with non-uniform meshes for nonlinear fractional differential equations [☆]



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ABSTRACT

In this article, finite difference methods with non-uniform meshes for solving nonlinear fractional differential equations are presented, where the non-equidistant stepsize is non-decreasing. The rectangle formula and trapezoid formula are proposed based on the non-uniform meshes. Combining the above two methods, we then establish the predictor–corrector scheme. The error and stability analysis are carefully investigated. At last, numerical examples are carried out to verify the theoretical analysis. Besides, the comparisons between non-uniform and uniform meshes are given, where the non-uniform meshes show the better performance when dealing with the less smooth problems.

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1. Introduction

In recent years, growing attention has been focused on fractional differential equations because they can provide a better approach to describe the complex phenomena in nature, such as viscoelastic materials, anomalous diffusion, signal processing and control theory, etc., see [1,16,17,28,29,32,38]. Compared to classical integer-order differential equations, the theoretical investigations and establishment of numerical schemes for fractional-order (or fractional for brevity) versions are more complicated due to the special properties of fractional differential operators, such as the non-locality, history dependence, and/or long-range interactions [37].

In this paper, we study the numerical schemes for the following nonlinear fractional initial value problem

$$\begin{cases} {}_c D_{0,t}^\alpha y(t) = f(t, y), & T \geq t > 0, \\ y^{(k)}(0) = y_0^{(k)}, & k = 0, 1, \dots, n-1, \end{cases} \quad (1.1)$$

where the continuous function $f(t, y)$ is nonlinear with respect to the unknown function y , n is a positive integer such that $n-1 < \alpha < n$ and the initial values $y_0^{(k)}$ are assumed to be given. Here ${}_c D_{0,t}^\alpha y(t)$ is the Caputo derivative, defined by

$${}_c D_{0,t}^\alpha y(t) = \frac{1}{\Gamma(n-\alpha)} \int_0^t (t-s)^{n-\alpha-1} y^{(n)}(s) ds, \quad n-1 < \alpha < n \in \mathbb{Z}^+.$$

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In order to ensure the existence of a unique solution of (1.1), we always assume that f satisfies Lipschitz condition with respect to the second variable, that is, $|f(t, y_1) - f(t, y_2)| \leq L|y_1 - y_2|$ where $L > 0$. It is well known that the initial value problem (1.1) is equivalent to the following Volterra integral equation [6,16]

$$y(t) = \sum_{j=0}^{n-1} \frac{t^j}{j!} y_0^{(j)} + \frac{1}{\Gamma(\alpha)} \int_0^t (t-s)^{\alpha-1} f(s, y(s)) ds, \quad (1.2)$$

in the sense that a continuous function is a solution of (1.1) if and only if it solves equation (1.2).

There have existed some studies on the numerical approaches for fractional differential equation (1.1). Early in 1986, Lubich studied the numerical approximation of fractional integrals based on a discrete convolution form and introduced the fractional linear multistep method [26,27]. Almost at the same time, Brunner and Houwen also presented linear multistep methods for Volterra equations [2]. In 1997, Diethelm proposed an implicit algorithm for the fractional differential equations based on a quadrature formula approach in a finite-part integral sense [5]. For more details about the relationship between fractional derivative and finite-part integral, see [12,13]. In 2000, Podlubny proposed a unification discrete form based on triangular strip matrices [33], which was further developed in [34]. In 2002, Diethelm et al. discussed an Adams-type predictor–corrector method for equation (1.2) and gave a detailed error analysis that the convergent order was $\min\{2, 1 + \alpha\}$ if ${}_C D_{0,t}^\alpha y(t) \in C^2[0, T]$ [7,8]. Shortly after, Li and Tao further studied the error analysis for the fractional predictor–corrector method [20]. And in [22], Li and Zeng reviewed the finite difference methods for fractional ordinary/partial differential equations. In [24], various kinds of numerical methods for fractional differential equations have been studied thoroughly. There have also been some investigations on the stability analysis for nonlinear fractional differential equations. Ladaci and Moulayin [18] analyzed the L^p -stability of fractional nonlinear differential equations. Deng [4] considered the sufficient conditions for the local asymptotical stability of nonlinear fractional differential equations. After that, Li and Zeng [23] studied the numerical stability of the finite difference methods for nonlinear fractional ordinary differential equation (1.1). In very recent, a series of numerical methods for Caputo derivatives is derived and the established algorithms are applied to Caputo-type advection–diffusion equations [3,19,21]. In the meantime, a series of numerical methods for Riesz derivatives is established and applied to Riesz-type fractional partial differential equations [9–11]. Keshtkar et al. [15] investigated the stability of equilibria in the nonlinear fractional-order dynamical systems.

In general, the existence of a weakly singular kernel $(t-s)^{\alpha-1}$ ($0 < \alpha < 1$) in fractional derivative and integral makes it more difficult to get a higher-order scheme. Particularly when the solution of equation (1.1) is not suitably smooth, those methods on uniform meshes seem to have a poor convergent rate. For these reasons, numerical methods on non-uniform meshes have been placed on the agenda. Especially in recent years, finite difference schemes with non-uniform meshes for fractional differential equations have attracted increasing attention. In [36,40], Yuste and Quintana-Murillo proposed an L1 method with non-uniform timesteps for fractional diffusion and diffusion–wave equations. In [35], Podlubny et al. studied the matrix approach on non-equidistant grids. Mustapha et al. [30,31] used the finite difference method to a sub-diffusion equation. Lopez-Fernandez and Sauter [25] presented a generalized convolution quadrature with variable timesteps. In very recent, Zhang et al. [42] investigated the finite difference scheme for the fractional diffusion equation on non-uniform meshes. Finite difference methods with non-uniform meshes often show great advantages when dealing with less smooth problems. However, the theoretical analysis of stability and convergence of the schemes with non-uniform meshes for nonlinear fractional differential equation seems not to be derived thoroughly yet.

In this paper, we mainly focus on the stability and convergence analysis of three kinds of numerical approaches on non-uniform meshes and illustrate their suitability for non-smooth problems through numerical tests. The rest of the paper is organized as follows. In Section 2, we outline the numerical schemes on the non-uniform meshes. Detailed stability and error analysis for the derived schemes are given in Sections 3 and 4, respectively. In Section 5, numerical examples are carried out to verify the theoretical analysis and to check the capability of the derived methods for non-smooth problems. The conclusions are included in the last section.

2. Numerical schemes on non-uniform meshes

For an integer N and the given time T , we divide the interval $[0, T]$ into $0 = t_0 < t_1 < \dots < t_k < t_{k+1} < \dots < t_N = T$, with non-equidistant stepsizes $\tau_i = t_{i+1} - t_i$, $0 \leq i \leq N-1$, and denote $\tau_{\max} = \max_{0 \leq i \leq N-1} \tau_i$, $\tau_{\min} = \min_{0 \leq i \leq N-1} \tau_i$. If the given question is singular at the origin, then the choice of the non-equidistant stepsizes obeys non-decreasing rule, i.e., $\tau_i \leq \tau_{i+1}$. Throughout this paper, we do not re-state this if no confusion appears.

Let y_j be the approximate solution of $y(t_j)$ ($j = 0, 1, \dots, k$) which have been determined. Now we need to calculate y_{k+1} . We introduce three approaches to do this.

Consider the following integral

$$\begin{aligned} I_{k+1} &= \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} f(s, y(s)) ds \\ &= \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} f(s, y(s)) ds, \end{aligned}$$

where $k = 0, 1, \dots, N - 1$.

It can be approximated by the following approach

$$I_{k+1} \approx \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} \tilde{f}_j(s, y(s)) ds,$$

where $\tilde{f}_j(s, y(s))$, $j = 0, 1, \dots, k$ is the approximation of $f(s, y(s))$ on the interval $[t_j, t_{j+1})$.

It will lead to different schemes by choosing different $\tilde{f}_j(s, y(s))$. Here we choose three kinds of $\tilde{f}_j(s, y(s))$ to derive the fractional rectangle, trapezoid, and predictor–corrector methods respectively.

(i) By choosing $\tilde{f}_j(s, y(s))$ as

$$\tilde{f}_j(s, y(s)) = f(t_j, y_j), \quad j = 0, 1, \dots, k,$$

the fractional rectangle method is derived as

$$y_{k+1} = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} y_0^{(j)} + \sum_{j=0}^k w_{j,k+1} f(t_j, y_j), \quad (2.1)$$

where

$$\begin{aligned} w_{j,k+1} &= \frac{1}{\Gamma(\alpha)} \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} ds \\ &= \frac{(t_{k+1} - t_j)^\alpha - (t_{k+1} - t_{j+1})^\alpha}{\Gamma(\alpha + 1)}, \quad j = 0, 1, \dots, k. \end{aligned} \quad (2.2)$$

(ii) If $\tilde{f}_j(s, y(s))$ is selected as

$$\tilde{f}_j(s, y(s)) = \frac{s - t_{j+1}}{t_j - t_{j+1}} f(t_j, y_j) + \frac{s - t_j}{t_{j+1} - t_j} f(t_{j+1}, y_{j+1}), \quad j = 0, 1, \dots, k,$$

then the fractional trapezoid method is given by

$$y_{k+1} = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} y_0^{(j)} + \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} f(t_j, y_j), \quad (2.3)$$

in which

$$\tilde{w}_{j,k+1} = \frac{1}{\Gamma(\alpha + 2)} \begin{cases} \frac{1}{t_1} A_0, & \text{if } j = 0, \\ \frac{1}{t_{j+1} - t_j} A_j + \frac{1}{t_{j-1} - t_j} B_j, & \text{if } j = 1, 2, \dots, k, \\ (t_{k+1} - t_k)^\alpha, & \text{if } j = k + 1, \end{cases} \quad (2.4)$$

and

$$\begin{cases} A_0 = (t_{k+1} - t_1)^{\alpha+1} - t_{k+1}^{\alpha+1} + (\alpha + 1)t_1 t_{k+1}^\alpha, \\ A_j = (t_{k+1} - t_{j+1})^{\alpha+1} - (t_{k+1} - t_j)^{\alpha+1} + (\alpha + 1)(t_{j+1} - t_j)(t_{k+1} - t_j)^\alpha, \\ B_j = (t_{k+1} - t_j)^{\alpha+1} - (t_{k+1} - t_{j-1})^{\alpha+1} + (\alpha + 1)(t_j - t_{j-1})(t_{k+1} - t_j)^\alpha. \end{cases}$$

After the observation of (2.3), it is obvious to see that this is an implicit scheme. In order to decrease computational complexity, predictor–corrector method is naturally proposed.

The predictor–corrector method can be deduced as the following steps. Firstly, we take equation (2.1) as the predictor item y_{k+1}^p and then we replace y_{k+1} on the right-hand side of scheme (2.3) with y_{k+1}^p to get the corrector item y_{k+1} which contributes to the scheme below.

(iii) The predictor–corrector method can be written as

$$\begin{cases} y_{k+1}^p = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} y_0^{(j)} + \sum_{j=0}^k w_{j,k+1} f(t_j, y_j), \\ y_{k+1} = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} y_0^{(j)} + \left\{ \sum_{j=0}^k \tilde{w}_{j,k+1} f(t_j, y_j) + \tilde{w}_{k+1,k+1} f(t_{k+1}, y_{k+1}^p) \right\}, \end{cases} \tag{2.5}$$

where $w_{j,k+1}$, $\tilde{w}_{j,k+1}$ are respectively defined by (2.2) and (2.4).

In order to obtain further higher-order numerical schemes on non-uniform meshes, we can also apply other quadrature rules to (1.2), such as Simpson rule, composite trapezoid rule and so on. We omit them here.

3. Stability analysis

In this section, we study the stability analysis for rectangle scheme (2.1), trapezoid scheme (2.3), and predictor–corrector scheme (2.5). Here the stability mainly refers to that if there is a perturbation in the initial condition, then the small change does not cause the large error in the numerical solution [23].

Suppose that y_k and z_k ($k = 1, 2, \dots, N$) are two solutions of the rectangle scheme (2.1) with different initial values $y_0^{(i)}$ and $z_0^{(i)}$ ($i = 0, 1, \dots, n - 1$), respectively. If there exists a positive constant C independent of non-equivalent stepsizes τ_j ($j = 0, 1, \dots, k$) and k , such that

$$|y_k - z_k| \leq C \sum_{i=0}^{n-1} |y_0^{(i)} - z_0^{(i)}|, \quad k = 1, 2, \dots, N,$$

then we say that rectangle scheme (2.1) is stable. It is similar to define the numerical stability for trapezoid scheme (2.3) and predictor–corrector scheme (2.5). Denote C as a generic positive constant that does not depend on meshes or k but on T, α and the smoothness of f if no ambiguousness occurs.

3.1. Several lemmas

Firstly, we introduce some lemmas that will be used in stability analysis as well as error analysis.

Lemma 3.1. *If $\alpha > 0$, k is a nonnegative integer, $\tau_j \leq \tau_{j+1}$ ($j = 0, 1, \dots, k - 1$), then $w_{j,k+1}$ and $\tilde{w}_{j,k+1}$ defined by equations (2.2) and (2.4) respectively have the following estimates*

$$w_{j,k+1} \leq C_\alpha \tau_j (t_{k+1} - t_j)^{\alpha-1}, \quad j = 0, 1, \dots, k, \tag{3.1}$$

and

$$\tilde{w}_{j,k+1} \leq C_\alpha \begin{cases} \tau_0 t_{k+1}^{\alpha-1}, & \text{if } j = 0, \\ [\tau_j (t_{k+1} - t_j)^{\alpha-1} + \tau_{j-1} (t_{k+1} - t_{j-1})^{\alpha-1}], & \text{if } j = 1, 2, \dots, k + 1, \end{cases} \tag{3.2}$$

where $C_\alpha = \frac{\max\{2, \alpha\}}{\Gamma(\alpha+1)}$.

Proof. It is simple to verify that

$$w_{k,k+1} = \frac{\tau_k^\alpha}{\Gamma(\alpha + 1)}$$

For $j = 0, 1, \dots, k - 1$, according to the mean value theorem, there exists a constant $\xi \in (t_j, t_{j+1})$, such that

$$\begin{aligned} \frac{\Gamma(\alpha + 1)w_{j,k+1}}{(t_{k+1} - t_j)^{\alpha-1}} &= \frac{(t_{k+1} - t_j)^\alpha - (t_{k+1} - t_{j+1})^\alpha}{(t_{k+1} - t_j)^{\alpha-1}} \\ &= \frac{\alpha \tau_j (t_{k+1} - \xi)^{\alpha-1}}{(t_{k+1} - t_j)^{\alpha-1}} \end{aligned}$$

If $\alpha \geq 1$, then

$$\frac{\Gamma(\alpha + 1)w_{j,k+1}}{(t_{k+1} - t_j)^{\alpha-1}} \leq \frac{\alpha\tau_j(t_{k+1} - t_j)^{\alpha-1}}{(t_{k+1} - t_j)^{\alpha-1}} \leq \alpha\tau_j.$$

If $0 < \alpha < 1$, then

$$\begin{aligned} \frac{\Gamma(\alpha + 1)w_{j,k+1}}{(t_{k+1} - t_j)^{\alpha-1}} &= \alpha\tau_j \left(\frac{t_{k+1} - t_j}{t_{k+1} - \xi} \right)^{1-\alpha} \\ &\leq \alpha\tau_j \left(\frac{t_{k+1} - t_j}{t_{k+1} - t_{j+1}} \right)^{1-\alpha} \\ &= \alpha\tau_j \left(1 + \frac{\tau_j}{\tau_k + \tau_{k+1} + \dots + \tau_{j+1}} \right)^{1-\alpha} \\ &\leq 2^{1-\alpha} \alpha\tau_j \\ &\leq 2\tau_j. \end{aligned}$$

So inequality (3.1) is proved and inequality (3.2) can be proved in a similar way. All this completes the proof. \square

Lemma 3.2 (Gronwall inequality). (See [41].) Assume that $\{k_n\}$ and $\{p_n\}$ are nonnegative sequences, $g_0 \geq 0$, and the sequence $\{\phi_n\}$ satisfies

$$\begin{cases} \phi_0 \leq g_0, \\ \phi_n \leq g_0 + \sum_{j=0}^{n-1} p_j + \sum_{j=0}^{n-1} k_j \phi_j, \quad n \geq 1. \end{cases}$$

Then

$$\phi_n \leq (g_0 + \sum_{j=0}^{n-1} p_j) \exp\left(\sum_{j=0}^{n-1} k_j\right), \quad n \geq 1.$$

Next, we propose a modified Gronwall inequality which is useful to prove the stability and error estimates for the derived numerical methods.

Lemma 3.3. Assume that $\alpha, C_0, T > 0$ and $b_{j,k} = C_0\tau_j(t_k - t_j)^{\alpha-1}$ ($j = 0, 1, \dots, k-1$) for $0 = t_0 < t_1 < \dots < t_N = T$, $k = 1, 2, \dots, N$ where N is a positive integer and $\tau_j = t_{j+1} - t_j$. Let g_0 be positive and the sequence $\{\psi_k\}$ meet

$$\begin{cases} \psi_0 \leq g_0, \\ \psi_k \leq \sum_{j=0}^{k-1} b_{j,k} \psi_j + g_0, \end{cases} \quad (3.3)$$

then

$$\psi_k \leq Cg_0, \quad k = 1, 2, \dots, N. \quad (3.4)$$

Proof. If $\alpha \geq 1$, one gets $b_{j,k} \leq C_0T^{\alpha-1}\tau_j$, that is

$$\psi_k \leq C_0T^{\alpha-1} \sum_{j=0}^{k-1} \tau_j \psi_j + g_0.$$

It naturally leads to (3.4) by using Lemma 3.2.

Next, we study the situation $0 < \alpha < 1$.

Noticing that $\psi_0 \leq g_0$ and using the inequality (3.3) repeatedly, we have

$$\begin{aligned}
 \psi_k &\leq g_0 + \sum_{j_1=0}^{k-1} b_{j_1,k} \psi_{j_1} \leq g_0 + \sum_{j_1=0}^{k-1} b_{j_1,k} (g_0 + \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} \psi_{j_2}) \\
 &= g_0 + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} + \sum_{j_1=0}^{k-1} b_{j_1,k} \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} \psi_{j_2} \\
 &\leq g_0 + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} + \sum_{j_1=0}^{k-1} b_{j_1,k} \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} (g_0 + \sum_{j_3=0}^{j_2-1} b_{j_3,j_2} \psi_{j_3}) \\
 &= g_0 + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} + \sum_{j_1=0}^{k-1} b_{j_1,k} \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} \sum_{j_3=0}^{j_2-1} b_{j_3,j_2} \psi_{j_3} \\
 &\quad \vdots \\
 &\leq g_0 + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} \sum_{j_3=0}^{j_2-1} b_{j_3,j_2} \\
 &\quad + \cdots + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} \sum_{j_2=0}^{j_1-1} b_{j_2,j_1} \cdots \sum_{j_k=0}^{j_{k-1}-1} b_{j_k,j_{k-1}} \\
 &= g_0 + g_0 \sum_{j_1=0}^{k-1} b_{j_1,k} + g_0 \sum_{j_2=0}^{k-1} b_{j_2,k} \sum_{j_1=0}^{j_2-1} b_{j_1,j_2} + g_0 \sum_{j_3=0}^{k-1} b_{j_3,k} \sum_{j_2=0}^{j_3-1} b_{j_2,j_3} \sum_{j_1=0}^{j_2-1} b_{j_1,j_2} \\
 &\quad + \cdots + g_0 \sum_{j_k=0}^{k-1} b_{j_k,k} \sum_{j_{k-2}=0}^{j_{k-1}-1} b_{j_{k-2},j_{k-1}} \cdots \sum_{j_1=0}^{j_2-1} b_{j_1,j_2}.
 \end{aligned} \tag{3.5}$$

In the following, we show that the right-hand side of the last equality of (3.5) is bounded. Since $f(x) = x^\beta(t_{j_s} - x)^{\alpha-1}$ ($0 < x < t_{j_s}, \beta \geq 0$) is monotonically increasing, one has

$$\begin{aligned}
 \sum_{j_r=0}^{j_s-1} b_{j_r,j_s} t_{j_r}^\beta &\triangleq C_0 \sum_{j_r=0}^{j_s-1} \tau_{j_r} (t_{j_s} - t_{j_r})^{\alpha-1} t_{j_r}^\beta, \quad r < s \\
 &= C_0 \sum_{j_r=0}^{j_s-1} \tau_{j_r} f(t_{j_r}) \\
 &\leq C_0 \int_0^{t_{j_s}} f(t) dt \\
 &= C_0 t_{j_s}^{\alpha+\beta} B(\alpha, \beta + 1),
 \end{aligned}$$

where $B(\cdot, \cdot)$ denotes Euler beta function defined by

$$B(w, z) = \int_0^1 (1-t)^{w-1} t^{z-1} dt, \quad z, w > 0.$$

Hence, we have

$$\begin{aligned}
 F_{r,k} &= g_0 \sum_{j_r=0}^{k-1} b_{j_r,k} \sum_{j_{r-1}=0}^{j_r-1} b_{j_{r-1},j_r} \cdots \sum_{j_1=0}^{j_2-1} b_{j_1,j_2} \\
 &\leq g_0 C_0 \sum_{j_r=0}^{k-1} b_{j_r,k} \sum_{j_{r-1}=0}^{j_r-1} b_{j_{r-1},j_r} \cdots \sum_{j_2=0}^{j_3-1} b_{j_2,j_3} t_{j_2}^\alpha B(\alpha, 1) \\
 &\leq g_0 C_0^2 \sum_{j_r=0}^{k-1} b_{j_r,k} \sum_{j_{r-1}=0}^{j_r-1} b_{j_{r-1},j_r} \cdots \sum_{j_3=0}^{j_4-1} b_{j_3,j_4} t_{j_3}^{2\alpha} B(\alpha, 1) B(\alpha, \alpha + 1) \\
 &\quad \vdots \\
 &\leq g_0 C_0^r t_k^{r\alpha} \prod_{j=0}^{r-1} B(\alpha, j\alpha + 1) \\
 &\leq g_0 C_0^r T^{r\alpha} \prod_{j=0}^{r-1} B(\alpha, j\alpha + 1) \\
 &= g_0 C_0^r T^{r\alpha} \cdot \frac{\Gamma(\alpha)}{\Gamma(\alpha + 1)} \cdot \frac{\Gamma(\alpha)\Gamma(\alpha + 1)}{\Gamma(2\alpha + 1)} \cdots \\
 &\quad \frac{\Gamma(\alpha)\Gamma((r-2)\alpha + 1)}{\Gamma((r-1)\alpha + 1)} \cdot \frac{\Gamma(\alpha)\Gamma((r-1)\alpha + 1)}{\Gamma(r\alpha + 1)} \\
 &= g_0 \frac{C_0^r T^{r\alpha} (\Gamma(\alpha))^r}{\Gamma(r\alpha + 1)},
 \end{aligned} \tag{3.6}$$

where $r = 1, 2, \dots, k$.

Denote $a_r = \frac{C_0^r T^{r\alpha} (\Gamma(\alpha))^r}{\Gamma(r\alpha + 1)}$, $r = 1, 2, \dots, k$. Then

$$\frac{a_{r+1}}{a_r} = C_0 T^\alpha \frac{\Gamma(\alpha)\Gamma(r\alpha + 1)}{\Gamma((r+1)\alpha + 1)}.$$

According to Stirling's formula in terms of the Euler gamma function, for large z one gets

$$\Gamma(z) = \sqrt{2\pi} z^{z-\frac{1}{2}} e^{-z} e^{\frac{\theta_z}{12z}}, \quad \theta_z \in (0, 1).$$

So, for large r one has

$$\begin{aligned}
 \frac{\Gamma(r\alpha + 1)}{\Gamma((r+1)\alpha + 1)} &= \frac{r\Gamma(r\alpha)}{(r+1)\Gamma((r+1)\alpha)} \\
 &= \left(\frac{r}{r+1}\right)^{r\alpha + \frac{1}{2}} \left(\frac{e}{\alpha}\right)^\alpha (r+1)^{-\alpha} \exp\left(\frac{\theta_r}{12r\alpha} - \frac{\theta_{r+1}}{12(r+1)\alpha}\right) \\
 &\leq C_1 (r+1)^{-\alpha},
 \end{aligned}$$

where $\theta_r, \theta_{r+1} \in (0, 1)$, $C_1 = \left(\frac{e}{\alpha}\right)^\alpha e^{\frac{1}{12\alpha}}$.

Thus,

$$0 \leq \lim_{r \rightarrow \infty} \frac{a_{r+1}}{a_r} \leq \lim_{r \rightarrow \infty} C_0 C_1 T^\alpha \Gamma(\alpha) (r+1)^{-\alpha} = 0,$$

i.e.,

$$\lim_{r \rightarrow \infty} \frac{a_{r+1}}{a_r} = 0,$$

which means that $\sum_{r=1}^{\infty} a_r$ is convergent. Hence, the right-hand side of the last equality of (3.5) is bounded, that is

$$\psi_k \leq g_0 + g_0 \sum_{r=1}^k a_r \leq C g_0, \tag{3.7}$$

which shows that (3.4) holds for the case $0 < \alpha < 1$. It follows that (3.4) is true for all $\alpha > 0$. \square

The above lemma is very useful for numerical stability analysis and error estimate. There have existed some other kinds of Gronwall inequalities, for example, see [14,39,41].

3.2. Stability analysis

In this section, we first give the stability analysis of rectangle scheme (2.1) in the following theorem.

Theorem 3.1. *Suppose that y_j ($j = 1, 2, \dots, k$) are the solutions of the rectangle scheme (2.1) where the non-equidistant stepsize is non-decreasing, $f(t, y)$ satisfies the Lipschitz condition with respect to the second argument y with a Lipschitz constant L on the existed interval of its unique solution. Then the rectangle scheme (2.1) is stable.*

Proof. Suppose $\tilde{y}_0^{(i)}$ ($i = 0, 1, \dots, n - 1$) and \tilde{y}_j ($j = 0, 1, \dots, k + 1$) are the perturbations of $y_0^{(i)}$ and y_j , respectively. It follows from (2.1) that

$$\begin{cases} y_{k+1} = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} y_0^{(j)} + \sum_{j=0}^k w_{j,k+1} f(t_j, y_j), \\ y_{k+1} + \tilde{y}_{k+1} = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} (y_0^{(j)} + \tilde{y}_0^{(j)}) + \sum_{j=0}^k w_{j,k+1} f(t_j, y_j + \tilde{y}_j). \end{cases}$$

So we obtain the following perturbation equation

$$\tilde{y}_{k+1} = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} \tilde{y}_0^{(j)} + \sum_{j=0}^k w_{j,k+1} (f(t_j, y_j + \tilde{y}_j) - f(t_j, y_j)).$$

Denoting by $\eta_0 = \max_{0 \leq j \leq n-1} |\tilde{y}_0^{(j)}|$, we get

$$\begin{aligned} |\tilde{y}_{k+1}| &= \left| \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} \tilde{y}_0^{(j)} + \sum_{j=0}^k w_{j,k+1} (f(t_j, y_j + \tilde{y}_j) - f(t_j, y_j)) \right| \\ &\leq \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} \eta_0 + L \sum_{j=0}^k w_{j,k+1} |\tilde{y}_j|. \end{aligned}$$

Using Lemmas 3.1 and 3.3, we reach that

$$\begin{aligned} |\tilde{y}_{k+1}| &\leq C_1 \left(\sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} \eta_0 \right) \\ &\leq C_1 \left(\sum_{j=0}^{n-1} \frac{T^j}{j!} \right) \eta_0 \\ &= C \eta_0, \end{aligned}$$

where $C = C_1 \left(\sum_{j=0}^{n-1} \frac{T^j}{j!} \right)$. The proof is thus ended. \square

Similar to the proof of Theorem 3.1, we have the following theorem whose proof is omitted here or left to the interested readers.

Theorem 3.2. *Suppose that y_j ($j = 1, 2, \dots, k$) are the solutions of the trapezoid scheme (2.3) where the non-equidistant stepsize is non-decreasing, $f(t, y)$ satisfies the Lipschitz condition with respect to the second argument y with a Lipschitz constant L on the existed interval of its unique solution. Then the trapezoid scheme (2.3) is stable.*

Next, we present the stability of scheme (2.5).

Theorem 3.3. *Suppose that y_j ($j = 1, 2, \dots, k$) are the solutions of the predictor–corrector scheme (2.5) where the non-equidistant stepsize is non-decreasing, $f(t, y)$ satisfies the Lipschitz condition with respect to the second argument y with a Lipschitz constant L on the existed interval of its unique solution. Then the predictor–corrector scheme (2.5) is stable.*

Proof. Suppose $\tilde{y}_0^{(i)}$ ($i = 0, 1, \dots, n-1$), \tilde{y}_j ($j = 0, 1, \dots, k+1$) and \tilde{y}_{k+1}^p ($k = 0, 1, \dots, N-1$) are the perturbations of $y_0^{(i)}$, y_j and y_{k+1}^p , respectively. Then we get the following perturbation equations

$$\begin{cases} \tilde{y}_{k+1}^p = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} \tilde{y}_0^{(j)} + \sum_{j=0}^k w_{j,k+1} \left(f(t_j, y_j + \tilde{y}_j) - f(t_j, y_j) \right), \\ \tilde{y}_{k+1} = \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} \tilde{y}_0^{(j)} + \left\{ \sum_{j=0}^k \tilde{w}_{j,k+1} \left(f(t_j, y_j + \tilde{y}_j) - f(t_j, y_j) \right) \right. \\ \left. + \tilde{w}_{k+1,k+1} \left(f(t_{k+1}, y_{k+1}^p + \tilde{y}_{k+1}^p) - f(t_{k+1}, y_{k+1}^p) \right) \right\}. \end{cases}$$

Denote by $\eta_0 = \max_{0 \leq j \leq n-1} |\tilde{y}_0^{(j)}|$. Noticing that the following estimate holds

$$\tilde{w}_{k+1,k+1} \leq \frac{T^\alpha}{\Gamma(\alpha + 2)}, \quad (3.8)$$

one has

$$\begin{aligned} |\tilde{y}_{k+1}| &= \left| \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} \tilde{y}_0^{(j)} + \left\{ \sum_{j=0}^k \tilde{w}_{j,k+1} \left(f(t_j, y_j + \tilde{y}_j) - f(t_j, y_j) \right) \right. \right. \\ &\quad \left. \left. + \tilde{w}_{k+1,k+1} \left(f(t_{k+1}, y_{k+1}^p + \tilde{y}_{k+1}^p) - f(t_{k+1}, y_{k+1}^p) \right) \right\} \right| \\ &\leq \sum_{j=0}^{n-1} \frac{t_{k+1}^j}{j!} |\tilde{y}_0^{(j)}| + L \sum_{j=0}^k \tilde{w}_{j,k+1} |\tilde{y}_j| + L \tilde{w}_{k+1,k+1} |\tilde{y}_{k+1}^p| \\ &\leq \eta_0 \sum_{j=0}^{n-1} \frac{T^j}{j!} + L \sum_{j=0}^k \tilde{w}_{j,k+1} |\tilde{y}_j| + \frac{LT^\alpha}{\Gamma(\alpha + 2)} |\tilde{y}_{k+1}^p| \\ &\leq \left(1 + \frac{LT^\alpha}{\Gamma(2 + \alpha)} \right) \eta_0 \sum_{j=0}^{n-1} \frac{T^j}{j!} + L \sum_{j=0}^k \left(\tilde{w}_{j,k+1} + \frac{LT^\alpha}{\Gamma(\alpha + 2)} w_{j,k+1} \right) |\tilde{y}_j|. \end{aligned}$$

According to [Lemmas 3.1 and 3.3](#), it leads to

$$\begin{aligned} |\tilde{y}_{k+1}| &\leq \left(1 + \frac{LT^\alpha}{\Gamma(2 + \alpha)} \right) \eta_0 \sum_{j=0}^{n-1} \frac{T^j}{j!} + LC_\alpha \left(2 + \frac{LT^\alpha}{\Gamma(\alpha + 2)} \right) \sum_{j=0}^k \tau_j (t_{k+1} - t_j)^{\alpha-1} |\tilde{y}_j| \\ &\leq C_1 \left(1 + \frac{LT^\alpha}{\Gamma(2 + \alpha)} \right) \eta_0 \sum_{j=0}^{n-1} \frac{T^j}{j!} \\ &= C\eta_0, \end{aligned}$$

where $C = C_1 \left(1 + \frac{LT^\alpha}{\Gamma(2 + \alpha)} \right) \sum_{j=0}^{n-1} \frac{T^j}{j!}$. All this ends the proof. \square

Remark 3.1. The global Lipschitz conditions in [Theorems 3.1–3.3](#) can be weakened to the local Lipschitz conditions. The proof techniques are almost the same as those studied in [\[23\]](#).

4. Error analysis

In this section, we give the error analysis of rectangle scheme [\(2.1\)](#), trapezoid scheme [\(2.3\)](#), and predictor–corrector scheme [\(2.5\)](#).

At first, we present some lemmas that will be used later on.

Lemma 4.1. If $g(t) \in C^1[0, T]$, then

$$\left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^k w_{j,k+1} g(t_j) \right| \leq \frac{\|g'\|_\infty}{\Gamma(\alpha + 1)} T^\alpha \tau_{max}.$$

Proof.

$$\begin{aligned} & \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^k w_{j,k+1} g(t_j) \right| \\ &= \frac{1}{\Gamma(\alpha)} \left| \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} (g(t_j) - g(s)) ds \right| \\ &\leq \frac{1}{\Gamma(\alpha)} \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} |(s - t_j) g'(\xi_j)| ds, \quad \xi_j \in (t_j, t_{j+1}) \\ &\leq \frac{1}{\Gamma(\alpha)} \|g'\|_\infty \tau_{max} \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} ds \\ &= \frac{1}{\Gamma(\alpha + 1)} \|g'\|_\infty t_{k+1}^\alpha \tau_{max} \leq \frac{\|g'\|_\infty}{\Gamma(\alpha + 1)} T^\alpha \tau_{max}. \end{aligned}$$

This finishes the proof. \square

Lemma 4.2. If $g(t) \in C^2[0, T]$, then

$$\left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} g(t_j) \right| \leq \frac{\|g''\|_\infty}{2\Gamma(\alpha + 1)} T^\alpha \tau_{max}^2.$$

Proof.

$$\begin{aligned} & \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} g(t_j) \right| \\ &= \frac{1}{\Gamma(\alpha)} \left| \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} \left\{ g(s) - \frac{s - t_{j+1}}{t_j - t_{j+1}} g(t_j) - \frac{s - t_j}{t_{j+1} - t_j} g(t_{j+1}) \right\} ds \right| \\ &\leq \frac{1}{\Gamma(\alpha)} \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} \left| \frac{1}{2} (s - t_{j+1})(s - t_j) g''(\xi_j) \right| ds, \quad \xi_j \in (t_j, t_{j+1}) \\ &\leq \frac{\|g''\|_\infty}{2\Gamma(\alpha)} \tau_{max}^2 \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} ds \\ &= \frac{\|g''\|_\infty}{2\Gamma(\alpha + 1)} t_{k+1}^\alpha \tau_{max}^2 \leq \frac{\|g''\|_\infty}{2\Gamma(\alpha + 1)} T^\alpha \tau_{max}^2. \end{aligned}$$

The proof is thus completed. \square

In the above two lemmas, the smooth conditions of the integrand $g(t)$ are sufficient ones which are somewhat strong. In such situations, non-uniform meshes seem not to have superiority. It means that the convergent rates of both non-uniform meshes and uniform ones are highly possible the same for smooth functions. But for non-smooth functions, non-uniform

meshes are much more suitable than uniform meshes. Now we give another lemma based on concrete non-uniform meshes to show this.

Lemma 4.3. Let $g(t) = t^\sigma$ ($0 < \sigma < 1$) which is not smooth at the origin $t = 0$. Set non-uniform meshes with variable stepsizes $\tau_j = t_{j+1} - t_j = (j+1)\mu$, $0 \leq j \leq N-1$, where $\mu = \frac{2T}{N(N+1)}$. Then for $0 < \alpha < 1$, we have

$$\begin{aligned} & \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} g(t_j) \right| \\ & \leq C_{\sigma,\alpha} N^{-2(\sigma+\alpha)} (k+1)^{2(\sigma+\alpha-1)}, \end{aligned} \quad (4.1)$$

where $k = 0, 1, \dots, N-1$, and $C_{\sigma,\alpha}$ is a constant depending only on T, α, σ .

Proof.

$$\begin{aligned} & \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} g(t_j) \right| \\ & = \frac{1}{\Gamma(\alpha)} \left| \sum_{j=0}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} \left\{ g(s) - \frac{s - t_{j+1}}{t_j - t_{j+1}} g(t_j) - \frac{s - t_j}{t_{j+1} - t_j} g(t_{j+1}) \right\} ds \right| \\ & \leq \frac{1}{\Gamma(\alpha)} \sum_{j=1}^k \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} \left| \frac{1}{2} (s - t_{j+1})(s - t_j) \sigma(\sigma - 1) \xi_j^{\sigma-2} \right| ds \\ & \quad + \frac{1}{\Gamma(\alpha)} \int_0^{t_1} (t_{k+1} - s)^{\alpha-1} (s^\sigma - s t_1^{\sigma-1}) ds \\ & = I_1 + I_2, \end{aligned} \quad (4.2)$$

where $\xi_j \in (t_j, t_{j+1})$. It is easy to show that the above inequality is bounded by $C_{\sigma,\alpha} t_1^{\sigma+\alpha}$ when $k = 0$.

Next we focus on the case $k \geq 1$. For the first part I_1 on the right-hand side of (4.1), we have the following estimate.

$$\begin{aligned} I_1 & \leq \frac{1}{2\Gamma(\alpha)} \sigma(1-\sigma) \sum_{j=1}^k \tau_j^2 t_j^{\sigma-2} \int_{t_j}^{t_{j+1}} (t_{k+1} - s)^{\alpha-1} ds \\ & \leq \frac{1}{2\Gamma(\alpha)} \sigma(1-\sigma) \sum_{j=1}^k \tau_j^3 t_j^{\sigma-2} (t_{k+1} - t_{j+1})^{\alpha-1} \\ & \leq C_{\sigma,\alpha} \left(\frac{T}{N(N+1)} \right)^{\sigma+\alpha} \sum_{j=1}^k (j+1)^{\sigma+1} j^{\sigma-2} (k+j+3)^{\alpha-1} (k-j)^{\alpha-1} \\ & \leq C_{\sigma,\alpha} \left(\frac{T}{N(N+1)} \right)^{\sigma+\alpha} \sum_{j=1}^k \frac{j+1}{j} (j+1)^\sigma j^{\sigma-1} (k+j+3)^{\alpha-1} (k-j)^{\alpha-1} \\ & \leq C_{\sigma,\alpha} \left(\frac{T}{N(N+1)} \right)^{\sigma+\alpha} (k+1)^{\sigma+\alpha-1} \sum_{j=1}^k j^{\sigma-1} (k-j)^{\alpha-1} \\ & \leq C_{\sigma,\alpha} \left(\frac{T}{N(N+1)} \right)^{\sigma+\alpha} (k+1)^{\sigma+\alpha-1} \int_0^k s^{\sigma-1} (k-s)^{\alpha-1} ds \\ & \leq C_{\sigma,\alpha} \left(\frac{T}{N(N+1)} \right)^{\sigma+\alpha} (k(k+1))^{\sigma+\alpha-1}. \end{aligned}$$

For $0 < \sigma + \alpha < 1$, we get that

$$\begin{aligned}
 I_1 &\leq C_{\sigma,\alpha} N^{-2(\sigma+\alpha)} k^{2(\sigma+\alpha-1)} \\
 &\leq C_{\sigma,\alpha} N^{-2(\sigma+\alpha)} (k+1)^{2(\sigma+\alpha-1)} \left(1 - \frac{1}{k+1}\right)^{2(\sigma+\alpha-1)} \\
 &\leq 4C_{\sigma,\alpha} N^{-2(\sigma+\alpha)} (k+1)^{2(\sigma+\alpha-1)}.
 \end{aligned} \tag{4.3}$$

Or else $1 \leq \sigma + \alpha < 2$, then $I_1 \leq C_{\sigma,\alpha} N^{-2(\sigma+\alpha)} (k+1)^{2(\sigma+\alpha-1)}$. Thus we can reach that $I_1 \leq C_{\sigma,\alpha} N^{-2(\sigma+\alpha)} (k+1)^{2(\sigma+\alpha-1)}$ with a generic constant for arbitrary $0 < \alpha, \sigma < 1$.

For the second part I_2 , one has

$$\begin{aligned}
 I_2 &= \frac{1}{\Gamma(\alpha)} \int_0^{t_1} (t_{k+1} - s)^{\alpha-1} (s^\sigma - st_1^{\sigma-1}) ds \\
 &\leq \frac{1}{\Gamma(\alpha)} \left(\frac{1}{\sigma+1} - \frac{1}{2}\right) t_1^{\sigma+1} (t_{k+1} - t_1)^{\alpha-1} \\
 &\leq C_{\sigma,\alpha} \left(\frac{T}{N(N+1)}\right)^{\sigma+\alpha} (k(k+1))^{\alpha-1} \\
 &\leq 4C_{\sigma,\alpha} N^{-2(\sigma+\alpha)} (k+1)^{2(\alpha-1)},
 \end{aligned}$$

where the technique of deriving (4.3) is used in the last inequality.

So, combining I_1 with I_2 , we can obtain inequality (4.1). The proof is complete. \square

Remark 4.1. For uniform meshes with a uniform stepsize $h = \frac{T}{N}$ and $g(t) = t^\sigma$ ($0 < \sigma < 1$), it has been proved in [8] that

$$\left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} g(t_j) \right| \leq C_{\sigma,\alpha} t_{k+1}^{\alpha-1} h^{\sigma+1},$$

i.e.,

$$\left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} g(s) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} g(t_j) \right| \leq C_{\sigma,\alpha} N^{-(\sigma+\alpha)} (k+1)^{\alpha-1}. \tag{4.4}$$

In view of inequalities (4.1) and (4.4), we find that the order of convergence with non-uniform meshes are higher than that with uniform meshes.

From the above Lemmas 4.1 and 4.2, one can easily get the following theorems.

Theorem 4.1. If ${}_C D_{0,t}^\alpha y(t) \in C^1[0, T]$ and the non-equidistant stepsize is non-decreasing, then the rectangle scheme (2.1) for equation (1.2) has the following estimate,

$$|y_{k+1} - y(t_{k+1})| \leq C \tau_{max}, \quad k = 0, 1, \dots, N - 1. \tag{4.5}$$

Theorem 4.2. If ${}_C D_{0,t}^\alpha y(t) \in C^2[0, T]$ and the non-equidistant stepsize is non-decreasing, then the trapezoid scheme (2.3) for equation (1.2) has the following estimate,

$$|y_{k+1} - y(t_{k+1})| \leq C \tau_{max}^2, \quad k = 0, 1, \dots, N - 1. \tag{4.6}$$

Next, we derive the error estimate of scheme (2.5).

Theorem 4.3. If ${}_C D_{0,t}^\alpha y(t) \in C^2[0, T]$ and the non-equidistant stepsize is non-decreasing, then the predictor–corrector scheme (2.5) has the estimate below,

$$|y_{k+1} - y(t_{k+1})| \leq C \tau_{max}^q, \quad k = 0, 1, \dots, N - 1, \tag{4.7}$$

where $q = \min\{2, 1 + \alpha\}$.

Proof. Since $f(t, y(t)) = {}_C D_{0,t}^\alpha y(t) \in C^2[0, T]$ is bounded, then there exists a constant $M > 0$, such that $|f_t| \leq M$ and $|f_{tt}| \leq M$. We denote $e_{k+1} = |y(t_{k+1}) - y_{k+1}|$ for $k = 0, 1, \dots, N - 1$. By tedious calculations, one has

$$\begin{aligned}
e_{k+1} &= \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} f(s, y(s)) ds \right. \\
&\quad \left. - \left\{ \sum_{j=0}^k \tilde{w}_{j,k+1} f(t_j, y_j) + \tilde{w}_{k+1,k+1} f(t_{k+1}, y_{k+1}^P) \right\} \right| \\
&\leq \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} f(s, y(s)) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} f(t_j, y(t_j)) \right| \\
&\quad + \left| \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} f(t_j, y(t_j)) - \left\{ \sum_{j=0}^k \tilde{w}_{j,k+1} f(t_j, y_j) \right. \right. \\
&\quad \left. \left. + \tilde{w}_{k+1,k+1} f(t_{k+1}, y_{k+1}^P) \right\} \right| \\
&\leq \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1} - s)^{\alpha-1} f(s, y(s)) ds - \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} f(t_j, y(t_j)) \right| \\
&\quad + \left| \sum_{j=0}^k \tilde{w}_{j,k+1} [f(t_j, y(t_j)) - f(t_j, y_j)] \right| \\
&\quad + \left| \tilde{w}_{k+1,k+1} [f(t_{k+1}, y(t_{k+1})) - f(t_{k+1}, y_{k+1}^P)] \right| \\
&\triangleq I_1 + I_2 + I_3.
\end{aligned}$$

It follows from Lemma 4.2 that

$$I_1 \leq C_1 \tau_{max}^2,$$

where $C_1 = \frac{MT^\alpha}{2\Gamma(1+\alpha)}$.

Since f is presumed to satisfy the Lipschitz condition with a Lipschitz constant L , one has

$$\begin{aligned}
I_2 &= \left| \sum_{j=0}^k \tilde{w}_{j,k+1} \{f(t_j, y(t_j)) - f(t_j, y_j)\} \right| \\
&\leq \sum_{j=0}^k \tilde{w}_{j,k+1} |f(t_j, y(t_j)) - f(t_j, y_j)| \\
&\leq L \sum_{j=0}^k \tilde{w}_{j,k+1} |y(t_j) - y_j| \\
&= L \sum_{j=0}^k \tilde{w}_{j,k+1} e_j.
\end{aligned}$$

By the definition of $\tilde{w}_{k+1,k+1}$, we get

$$\tilde{w}_{k+1,k+1} = \frac{1}{\Gamma(\alpha+2)} (t_{k+1} - t_k)^\alpha = \frac{1}{\Gamma(\alpha+2)} \tau_k^\alpha \leq \frac{1}{\Gamma(\alpha+2)} \tau_{max}^\alpha.$$

Thus

$$\begin{aligned}
I_3 &= \left| \tilde{w}_{k+1,k+1} \{f(t_{k+1}, y(t_{k+1})) - f(t_{k+1}, y_{k+1}^P)\} \right| \\
&\leq \frac{1}{\Gamma(\alpha+2)} \tau_{max}^\alpha |f(t_{k+1}, y(t_{k+1})) - f(t_{k+1}, y_{k+1}^P)| \\
&\leq \frac{L}{\Gamma(\alpha+2)} \tau_{max}^\alpha |y(t_{k+1}) - y_{k+1}^P|
\end{aligned}$$

$$\begin{aligned}
 &= \frac{L\tau_{max}^\alpha}{\Gamma(\alpha+2)} \cdot \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1}-s)^{\alpha-1} f(s, y(s)) ds - \sum_{j=0}^k w_{j,k+1} f(t_j, y_j) \right| \\
 &\leq \frac{L\tau_{max}^\alpha}{\Gamma(\alpha+2)} \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1}-s)^{\alpha-1} {}_C D_{0,s}^\alpha y(s) ds - \sum_{j=0}^k w_{j,k+1} {}_C D_{0,t_j}^\alpha y(t_j) \right| \\
 &\quad + \frac{L\tau_{max}^\alpha}{\Gamma(\alpha+2)} \left| \sum_{j=0}^k w_{j,k+1} {}_C D_{0,t_j}^\alpha y(t_j) - \sum_{j=0}^k w_{j,k+1} f(t_j, y_j) \right| \\
 &\leq \frac{L\tau_{max}^\alpha}{\Gamma(\alpha+2)} \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_{k+1}} (t_{k+1}-s)^{\alpha-1} {}_C D_{0,s}^\alpha y(s) ds - \sum_{j=0}^k w_{j,k+1} {}_C D_{0,t_j}^\alpha y(t_j) \right| \\
 &\quad + \frac{L\tau_{max}^\alpha}{\Gamma(\alpha+2)} \sum_{j=0}^k w_{j,k+1} |f(t_j, y(t_j)) - f(t_j, y_j)| \\
 &\leq C_2 \tau_{max}^{\alpha+1} + \frac{L^2 T^\alpha}{\Gamma(\alpha+2)} \sum_{j=0}^k w_{j,k+1} e_j,
 \end{aligned}$$

where Lemma 4.1 is utilized and $C_2 = \frac{LMT^\alpha}{\Gamma(\alpha+1)\Gamma(\alpha+2)}$.

Therefore,

$$\begin{aligned}
 e_{k+1} &\leq I_1 + I_2 + I_3 \\
 &\leq C_1 \tau_{max}^2 + C_2 \tau_{max}^{\alpha+1} + L \sum_{j=0}^k \tilde{w}_{j,k+1} e_j + \frac{L^2 T^\alpha}{\Gamma(\alpha+2)} \sum_{j=0}^k w_{j,k+1} e_j \\
 &= C_1 \tau_{max}^2 + C_2 \tau_{max}^{\alpha+1} + L \sum_{j=0}^k \left\{ \tilde{w}_{j,k+1} + \frac{L T^\alpha}{\Gamma(\alpha+2)} w_{j,k+1} \right\} e_j.
 \end{aligned}$$

Then, it follows from Lemmas 3.1 and 3.3 again that

$$\begin{aligned}
 e_{k+1} &\leq C_3 (C_1 \tau_{max}^2 + C_2 \tau_{max}^{\alpha+1}) \\
 &\leq C \tau_{max}^q,
 \end{aligned}$$

where $q = \min\{2, 1 + \alpha\}$. The proof is thus complete. \square

Remark 4.2. The smooth condition ${}_C D_{0,t}^\alpha y(t) \in C^2[0, T]$ is a sufficient one. For smooth functions in $[0, T]$, the convergent rate of non-uniform meshes and uniform ones are almost the same. For non-smooth functions in $[0, T]$ (i.e. whose derivative(s) may not exist at $t = 0$ or other points), however, the non-uniform meshes behave better than the uniform meshes.

Remark 4.3. The conclusions in Theorems 4.1–4.3 still hold if the global Lipschitz conditions are reduced to the local Lipschitz ones [23].

5. Numerical examples

In this section, we present four numerical examples as follows.

If $\tau_j = \text{const.}$, $j = 0, 1, \dots, N - 1$, that is, equidistant division on the interval $[0, T]$, then the convergence rate of predictor–corrector scheme reduced to $2 - \alpha$ for $\alpha \in [1/2, 1)$ when $y(t)$ satisfies $y(t) \in C^2[0, T]$. This drawback will be overcome when the non-uniform meshes are used.

We adopt the non-uniform meshes defined as

$$\tau_k = (k + 1)\mu, \quad 0 \leq k \leq N - 1,$$

where $\mu = \frac{2T}{N(N+1)}$. In the following, we always use these non-uniform meshes to solve the fractional differential equations (see Examples 5.1–5.3).

Now, we give an example to test the error estimates.

Table 1
Absolute errors at $t = 1$ for equation (5.1) using rectangle scheme (2.1).

N	$\alpha = 0.25$	EOC	$\alpha = 0.75$	EOC	$\alpha = 1.25$	EOC	$\alpha = 1.75$	EOC
80	8.64E-02		6.40E-02		4.99E-02		3.63E-02	
160	3.95E-02	1.13	3.15E-02	1.02	2.53E-02	0.98	1.85E-02	0.97
320	1.85E-02	1.09	1.56E-02	1.02	1.28E-02	0.99	9.35E-03	0.99
640	8.85E-03	1.07	7.72E-03	1.01	6.41E-03	0.99	4.70E-03	0.99
1280	4.26E-03	1.05	3.84E-03	1.01	3.21E-03	1.00	2.36E-03	1.00
2560	2.06E-03	1.05	1.91E-03	1.00	1.61E-03	1.00	1.18E-03	1.00

Table 2
Absolute errors at $t = 1$ for equation (5.1) using trapezoid scheme (2.3).

N	$\alpha = 0.25$	EOC	$\alpha = 0.75$	EOC	$\alpha = 1.25$	EOC	$\alpha = 1.75$	EOC
80	7.95E-04		1.15E-03		8.96E-04		5.73E-04	
160	2.14E-04	1.89	2.90E-04	1.98	2.25E-04	1.99	1.44E-04	1.99
320	5.66E-05	1.92	7.30E-05	1.99	5.65E-05	2.00	3.62E-05	2.00
640	1.48E-05	1.94	1.83E-05	1.99	1.41E-05	2.00	9.05E-06	2.00
1280	3.83E-06	1.95	4.59E-06	2.00	3.54E-06	2.00	2.27E-06	2.00
2560	9.84E-07	1.96	1.15E-06	2.00	8.85E-07	2.00	5.67E-07	2.00

Example 5.1. Consider the following FODE

$${}_C D_{0,t}^\alpha y(t) = \frac{\Gamma(6)}{\Gamma(6-\alpha)} t^{5-\alpha} + \frac{36}{\Gamma(5-\alpha)} t^{4-\alpha} - \Gamma(3+\alpha)t^2 - y^2 + (t^5 + \frac{3}{2}t^4 - 2t^{2+\alpha})^2, \quad (5.1)$$

with the initial values

$$\begin{cases} y(0) = 0, y'(0) = 0, & \text{if } 1 < \alpha < 2, \\ y(0) = 0, & \text{if } 0 < \alpha < 1. \end{cases}$$

The exact solution of this equation is $y(t) = t^5 + \frac{3}{2}t^4 - 2t^{2+\alpha}$, and

$${}_C D_{0,t}^\alpha y(t) = \frac{\Gamma(6)}{\Gamma(6-\alpha)} t^{5-\alpha} + \frac{36}{\Gamma(5-\alpha)} t^{4-\alpha} - \Gamma(3+\alpha)t^2, \quad 0 < \alpha < 2.$$

Thus ${}_C D_{0,t}^\alpha y(t) \in C^2[0, T]$ for arbitrary $T > 0$ which fulfills the conditions of [Theorems 4.1–4.3](#). For convenience, we check only the rectangle and the trapezoid methods on non-uniform meshes. The predictor–corrector method on non-uniform meshes is omitted here or left to the interested readers. The numerical results are displayed in [Tables 1 and 2](#). According to [Theorems 4.1 and 4.3](#), the orders of convergence of these two methods should be 1 and 2, respectively. From [Tables 1 and 2](#) we can see that the experimental orders of convergence (EOC) support the theoretical convergent orders.

Example 5.2. Consider the following nonlinear equation

$${}_C D_{0,t}^\alpha y(t) = \begin{cases} \frac{2}{\Gamma(3-\alpha)} t^{2-\alpha} - y^2 + (t^2 - t + 1)^2, & \text{for } 1 < \alpha < 2, \\ \frac{2}{\Gamma(3-\alpha)} t^{2-\alpha} - \frac{1}{\Gamma(2-\alpha)} t^{1-\alpha} \\ - y^2 + (t^2 - t + 1)^2, & \text{for } 0 < \alpha < 1. \end{cases} \quad (5.2)$$

The initial values has been chosen as follows

$$\begin{cases} y(0) = 1, y'(0) = -1, & \text{if } 1 < \alpha < 2, \\ y(0) = 1, & \text{if } 0 < \alpha < 1. \end{cases}$$

The exact solution of equation (5.2) is $y(t) = t^2 - t + 1$. It is easy to verify that $y(t) \in C^2[0, T]$, but ${}_C D_{0,t}^\alpha y(t)$ doesn't belong to $C^2[0, T]$. We can observe from [Table 3](#) that the experimental order of convergence (EOC) on non-uniform meshes is $\min\{2, 1 + \alpha\}$ which is still in line with the theoretical analysis in [Theorem 4.3](#) and is even better than the theoretical order when $1 < \alpha < 2$.

The numerical results show that the method on non-uniform meshes has a much higher order of convergence than that on uniform meshes except the case when $\alpha = 0.3$. From this table, the EOCs of non-uniform and uniform meshes are almost the same for $\alpha = 0.3$. But when α increases from 0.3, the regularity of right-hand side of equation (5.2) becomes weaker

Table 3
Absolute errors at $t = 1$ for equation (5.2) using predictor–corrector scheme (2.5).

Meshes	N	$\alpha = 0.30$	EOC	$\alpha = 0.90$	EOC	$\alpha = 1.30$	EOC	$\alpha = 1.90$	EOC
Non-uniform	10	1.75E–01		2.20E–02		5.20E–03		8.43E–03	
	20	7.68E–02	1.19	5.93E–03	1.89	1.10E–03	2.24	2.11E–03	2.00
	40	2.72E–02	1.50	1.57E–03	1.92	2.23E–04	2.31	5.10E–04	2.05
	80	8.90E–03	1.61	4.14E–04	1.92	4.34E–05	2.36	1.22E–04	2.07
	160	2.95E–03	1.59	1.10E–04	1.92	8.14E–06	2.41	2.89E–05	2.07
	320	1.01E–03	1.54	2.91E–05	1.91	1.47E–06	2.47	6.88E–06	2.07
Uniform	10	8.57E–02		1.18E–02		1.50E–04		5.04E–02	
	20	2.86E–02	1.59	4.12E–03	1.52	1.97E–04	–0.39	2.37E–02	1.09
	40	9.07E–03	1.66	1.57E–03	1.39	1.12E–04	0.82	1.11E–02	1.10
	80	2.96E–03	1.61	6.47E–04	1.28	4.50E–05	1.31	5.18E–03	1.10
	160	1.01E–03	1.55	2.79E–04	1.21	1.61E–05	1.48	2.42E–03	1.10
	320	3.58E–04	1.50	1.24E–04	1.17	5.42E–06	1.57	1.13E–03	1.10

Table 4
Absolute errors at $t = 1$ for equation (5.3) using trapezoid scheme (2.3).

Meshes	N	$\alpha = 0.20$	EOC	$\alpha = 0.40$	EOC	$\alpha = 0.60$	EOC	$\alpha = 0.80$	EOC
Non-uniform	10	3.41E–04		3.69E–04		5.16E–04		6.24E–04	
	20	6.94E–05	2.30	9.08E–05	2.02	1.32E–04	1.96	1.62E–04	1.94
	40	1.49E–05	2.22	2.25E–05	2.01	3.36E–05	1.98	4.15E–05	1.97
	80	3.37E–06	2.15	5.59E–06	2.01	8.49E–06	1.99	1.05E–05	1.98
	160	7.85E–07	2.10	1.39E–06	2.00	2.13E–06	1.99	2.64E–06	1.99
	320	1.87E–07	2.07	3.48E–07	2.00	5.35E–07	2.00	6.63E–07	2.00
Uniform	10	1.29E–03		1.27E–03		9.61E–04		6.37E–04	
	20	5.46E–04	1.24	4.67E–04	1.44	3.07E–04	1.65	1.79E–04	1.84
	40	2.34E–04	1.22	1.74E–04	1.43	9.86E–05	1.64	5.03E–05	1.83
	80	1.01E–04	1.21	6.51E–05	1.42	3.19E–05	1.63	1.42E–05	1.82
	160	4.39E–05	1.21	2.45E–05	1.41	1.04E–05	1.62	4.01E–06	1.82
	320	1.91E–05	1.20	9.24E–06	1.41	3.39E–06	1.62	1.14E–06	1.82

due to the factors $t^{2-\alpha}$ and $t^{1-\alpha}$. In this situation, the convergent order of non-uniform meshes is obviously better than that of uniform meshes.

Example 5.3. Consider the following fractional differential equation

$${}_C D_{0,t}^\alpha y(t) + y(t) = 0, \quad y(0) = 1, \quad 0 < \alpha < 1. \tag{5.3}$$

The analytical solution is

$$y(t) = E_{\alpha,1}(-t^\alpha),$$

where $E_{\alpha,\beta}(z)$ is the Mittag–Leffler function defined by

$$E_{\alpha,\beta}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)}, \quad \alpha, \beta > 0.$$

Here we compare the trapezoid method (2.3) on uniform meshes with that on non-uniform meshes. And we choose the same number n of iteration, for example, $n = 10$ for both methods. From Table 4, one can see that the convergent order of this method on non-uniform meshes is still the second order though $y(t)$ is not smooth at $t = 0$, and that the present method on non-uniform meshes has much better performance than that on uniform meshes.

In the following, we give the last example.

Example 5.4. Consider the following time-fractional subdiffusion equation

$$\begin{cases} {}_C D_{0,t}^\alpha u(x, t) = \Delta u + f(x, t), & (x, t) \in (a, b) \times [0, T], \\ u(x, 0) = \phi(x), & x \in (a, b), \\ u(a, t) = \psi_1(t), \quad u(b, t) = \psi_2(t), & t \in [0, T]. \end{cases} \tag{5.4}$$

where Δ is Laplacian and $\alpha \in (0, 1)$.

Table 5
Errors in maximum norm at $t = 1$ for equation (5.4) using trapezoid scheme (5.7).

Mesher	N	$\alpha = 0.25$	EOC	$\alpha = 0.4$	EOC	$\alpha = 0.6$	EOC	$\alpha = 0.75$	EOC
Non-uniform	10	5.96E-05		7.43E-05		6.16E-05		3.28E-05	
	20	1.60E-05	1.89	1.94E-05	1.94	1.57E-05	1.98	8.31E-06	1.98
	40	4.17E-06	1.94	4.92E-06	1.98	3.87E-06	2.02	2.01E-06	2.05
	80	1.02E-06	2.03	1.19E-06	2.05	8.98E-07	2.11	4.28E-07	2.23
	160	1.96E-07	2.38	2.33E-07	2.35	1.55E-07	2.53	3.63E-08	3.56
Uniform	10	6.03E-05		9.65E-05		1.41E-04		1.63E-04	
	20	1.74E-05	1.79	2.93E-05	1.72	4.95E-05	1.51	6.62E-05	1.30
	40	5.01E-06	1.80	9.06E-06	1.69	1.82E-05	1.44	2.85E-05	1.22
	80	1.41E-06	1.83	2.83E-06	1.68	6.97E-06	1.39	1.27E-05	1.16
	160	3.55E-07	1.99	8.69E-07	1.71	2.73E-06	1.35	5.83E-06	1.13

Here we use the centered difference scheme to numerically approximate Δu in (5.4) with the spatial stepsize h , then the following fractional ordinary differential equation can be obtained

$${}_C D_{0,t}^\alpha u(x,t) = \frac{1}{h^2}(u(x+h,t) - 2u(x,t) + u(x-h,t)) + f(x,t). \quad (5.5)$$

Denote $x_i = a + ih$, $i = 0, 1, \dots, L$, and $h = \frac{b-a}{L}$. According to the trapezoid approach (2.3) for equation (1.1), we can derive that for equation (5.5) as follows with the giving initial-boundary condition.

$$u_i^{k+1} = \phi(x_i) + \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} f_i^j + \frac{1}{h^2} \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} (u_{i+1}^j - 2u_i^j + u_{i-1}^j), \quad (5.6)$$

$$1 \leq i \leq L-1, \quad 0 \leq k \leq N-1,$$

where u_i^j is the approximation of $u(x_i, t_j)$, $f_i^j = f(x_i, t_j)$ and $\tilde{w}_{j,k+1}$ is given by equation (2.4).

This implicit scheme can be rewritten in another form as below

$$\begin{aligned} & -\frac{1}{h^2} \tilde{w}_{k+1,k+1} u_{i+1}^{k+1} + \left(1 + \frac{2}{h^2} \tilde{w}_{k+1,k+1}\right) u_i^{k+1} - \frac{1}{h^2} \tilde{w}_{k+1,k+1} u_{i-1}^{k+1} \\ & = \phi(x_i) + \sum_{j=0}^{k+1} \tilde{w}_{j,k+1} f_i^j + \frac{1}{h^2} \sum_{j=0}^k \tilde{w}_{j,k+1} (u_{i+1}^j - 2u_i^j + u_{i-1}^j) \end{aligned} \quad (5.7)$$

where $i = 1, 2, \dots, L-1$ and $k = 0, 1, \dots, N-1$.

Now, we give the numerical experiment using Thomas algorithm based on scheme (5.7). If we take $u(x,t) = \sin(\pi x)(t^{0.8} + 1)$ as the solution of equation (5.4), then the source item and definite conditions are read as

$$\begin{cases} f(x,t) = \sin(\pi x) \left(\pi^2 (t^{0.8} + 1) + \frac{\Gamma(1.8)}{\Gamma(1.8 - \alpha)} t^{0.8 - \alpha} \right), & (x,t) \in (0,1) \times [0,T], \\ u(x,0) = \sin(\pi x), & x \in (0,1), \\ u(0,t) = 0, \quad u(1,t) = 0, & t \in [0,T]. \end{cases}$$

Taking the spatial stepsize as $h = \frac{1}{4096}$, the tests of convergent rate and absolute error on temporal direction are given below. The temporal non-uniform meshes are defined as

$$\tau_k = (k+1)\mu, \quad 0 \leq k \leq N-1,$$

where $\mu = \frac{2T}{N(N+1)}$ which are the same as those of Examples 5.1–5.3.

The numerical results are displayed in Table 5. Despite the nonsmoothness of the solution $u(x,t) = \sin(\pi x)(t^{0.8} + 1)$ on temporal direction at the initial time, it is easy to observe that the order of convergence on non-uniform meshes is still the second order which is much higher than that on uniform meshes.

6. Conclusion

In this article, we analyze the stability, convergence and error estimates for three kinds of fractional numerical methods on non-uniform meshes where the non-equidistant stepsize is non-decreasing. The numerical results show that all these methods on non-uniform meshes have better convergence and stability than those on uniform meshes both for FODEs and FPDEs with less smoothness. Particularly, the poorer smoothness of the solution and the right-hand side of equation (1.1), the greater advantages of the non-uniform meshes methods. We can naturally extend the methods and techniques used here to other numerical approaches on non-uniform meshes for fractional (ordinary/partial) differential equations with Caputo derivatives [3,19,21] and/or Riesz derivatives [9–11].

References

- [1] J.P. Bouchaud, A. Georges, Anomalous diffusion in disordered media: statistical mechanisms, models and physical applications, *Phys. Rep.* 195 (4) (1990) 127–293.
- [2] H. Brunner, P.J. van der Houwen, *The Numerical Solution of Volterra Equations*, North Holland, Amsterdam, 1986.
- [3] J.X. Cao, C.P. Li, Y.Q. Chen, High-order approximation to Caputo derivatives and Caputo-type advection–diffusion equations (II), *Fract. Calc. Appl. Anal.* 18 (2015) 735–761.
- [4] W.H. Deng, Smoothness and stability of the solutions for nonlinear fractional differential equations, *Nonlinear Anal.* 72 (3) (2010) 1768–1777.
- [5] K. Diethelm, An algorithm for the numerical solution of differential equations of fractional order, *Electron. Trans. Numer. Anal.* 5 (1) (1997) 1–6.
- [6] K. Diethelm, N.J. Ford, Analysis of fractional differential equations, *J. Math. Anal. Appl.* 265 (2) (2002) 229–248.
- [7] K. Diethelm, N.J. Ford, A.D. Freed, A predictor–corrector approach for the numerical solution of fractional differential equations, *Nonlinear Dyn.* 29 (2002) 3–22.
- [8] K. Diethelm, N.J. Ford, A.D. Freed, Detailed error analysis for a fractional Adams method, *Numer. Algorithms* 36 (1) (2004) 31–52.
- [9] H.F. Ding, C.P. Li, High-order algorithms for Riesz derivative and their applications (III), *Fract. Calc. Appl. Anal.* 19 (2016) 19–55.
- [10] H.F. Ding, C.P. Li, Y.Q. Chen, High-order algorithms for Riesz derivative and their applications (I), *Abstr. Appl. Anal.* (2014) 653797.
- [11] H.F. Ding, C.P. Li, Y.Q. Chen, High-order algorithms for Riesz derivative and their applications (II), *J. Comput. Phys.* 293 (2015) 218–237.
- [12] D. Elliott, An asymptotic analysis of two algorithms for certain Hadamard finite-part integrals, *IMA J. Numer. Anal.* 13 (3) (1993) 445–462.
- [13] D. Elliott, Three algorithms for Hadamard finite-part integrals and fractional derivatives, *J. Comput. Appl. Math.* 62 (3) (1995) 267–283.
- [14] Z.Q. Gong, D.L. Qian, C.P. Li, P. Guo, On the Hadamard type fractional differential system, in: D. Baleanu, J.A. Tenreiro Machado, Albert C.J. Luo (Eds.), *Fractional Dynamics and Control*, Springer, New York, 2012, pp. 159–171.
- [15] F. Keshtkar, G. Erjaee, M. Boutefnouchet, On stability of equilibrium points in nonlinear fractional differential equations and fractional Hamiltonian systems, *Complexity* 21 (2) (2015) 93–99.
- [16] A.A. Kilbas, H.M. Srivastava, J.J. Trujillo, *Theory and Application of Fractional Differential Equations*, Elsevier, Amsterdam, 2006.
- [17] R. Koeller, Applications of fractional calculus to the theory of viscoelasticity, *J. Appl. Mech.* 51 (2) (1984) 299–307.
- [18] S. Ladaci, E. Moulay, L^p -stability analysis of a class of nonlinear fractional differential equations, in: *International Conference on Electrical Engineering Design and Technologies*, hal-00177566, Hammamet, Tunisia, 2007, 4 pages.
- [19] H.F. Li, J.X. Cao, C.P. Li, High-order approximation to Caputo derivatives and Caputo-type advection–diffusion equations (III), *J. Comput. Appl. Math.* 299 (2016) 159–175, <http://www.mathworks.com/matlabcentral/fileexchange/53924>, *ibid* 53925.
- [20] C.P. Li, C.X. Tao, On the fractional Adams method, *Comput. Math. Appl.* 58 (8) (2009) 1573–1588.
- [21] C.P. Li, R.F. Wu, H.F. Ding, High-order approximation to Caputo derivatives and Caputo-type advection–diffusion equations, *Commun. Appl. Ind. Math.* 6 (2014) 1–32, e-536.
- [22] C.P. Li, F.H. Zeng, Finite difference methods for fractional differential equations, *Int. J. Bifurc. Chaos Appl. Sci. Eng.* 22 (4) (2012) 1230014.
- [23] C.P. Li, F.H. Zeng, The finite difference methods for fractional ordinary differential equations, *Numer. Funct. Anal. Optim.* 34 (2) (2013) 149–179.
- [24] C.P. Li, F.H. Zeng, *Numerical Methods for Fractional Calculus*, Chapman and Hall/CRC Press, Boca Raton, USA, 2015.
- [25] M. Lopez-Fernandez, S. Sauter, Generalized convolution quadrature with variable time stepping, *IMA J. Numer. Anal.* 33 (4) (2013) 1156–1175.
- [26] C. Lubich, Discretized fractional calculus, *SIAM J. Math. Anal.* 17 (3) (1986) 704–719.
- [27] C. Lubich, Convolution quadrature and discretized operational calculus, II, *Numer. Math.* 52 (4) (1988) 413–425.
- [28] R. Metzler, J. Klafter, The random walk’s guide to anomalous diffusion: a fractional dynamics approach, *Phys. Rep.* 339 (1) (2000) 1–77.
- [29] K. Miller, B. Ross, *An Introduction to the Fractional Calculus and Fractional Differential Equations*, Wiley, New York, 1993.
- [30] K. Mustapha, An implicit finite-difference time-stepping method for a sub-diffusion equation, with spatial discretization by finite elements, *IMA J. Numer. Anal.* 31 (2) (2011) 719–739.
- [31] K. Mustapha, J. AlMutawa, A finite difference method for an anomalous sub-diffusion equation, theory and applications, *Numer. Algorithms* 61 (4) (2012) 525–543.
- [32] I. Podlubny, *Fractional Differential Equations*, Academic Press, New York, 1999.
- [33] I. Podlubny, Matrix approach to discrete fractional calculus, *Fract. Calc. Appl. Anal.* 3 (4) (2000) 359–386.
- [34] I. Podlubny, A. Chechkin, T. Skovranek, Y. Chen, Blas M. Vinagre, Matrix approach to discrete fractional calculus, II: partial fractional differential equations, *J. Comput. Phys.* 228 (8) (2009) 3137–3153.
- [35] I. Podlubny, T. Skovranek, Blas M. Vinagre, I. Petras, V. Verbitsky, Y. Chen, Matrix approach to discrete fractional calculus, III: non-equidistant grids, variable step length and distributed orders, *Phil. Trans. R. Soc. A* 371 (1990) (2013) 20120153.
- [36] J. Quintana-Murillo, S.B. Yuste, A finite difference method with non-uniform timesteps for fractional diffusion and diffusion–wave equations, *Eur. Phys. J. Spec. Top.* 222 (8) (2013) 1987–1998.
- [37] S.G. Samko, A.A. Kilbas, O.I. Marichev, *Fractional Integrals and Derivatives Theory and Applications*, Gordon and Breach, New York, 1993.
- [38] H. Sheng, Y. Chen, T. Qiu, *Fractional Processes and Fractional-Order Signal Processing: Techniques and Applications*, Springer-Verlag, London, 2011.
- [39] H.P. Ye, J.M. Gao, Y.S. Ding, A generalized Gronwall inequality and its application to a fractional differential equation, *J. Math. Anal. Appl.* 328 (2007) 1075–1081.
- [40] S.B. Yuste, J. Quintana-Murillo, A finite difference method with non-uniform timesteps for fractional diffusion equations, *Comput. Phys. Commun.* 183 (12) (2012) 2594–2600.
- [41] F.H. Zeng, J.X. Cao, C.P. Li, Gronwall inequalities, in: C.P. Li, Y.J. Wu, R.S. Ye (Eds.), *Recent Advances in Applied Nonlinear Dynamics with Numerical Analysis*, World Scientific, Singapore, 2013, pp. 1–22.
- [42] Y.N. Zhang, Z.Z. Sun, H.L. Liao, Finite difference methods for the time fractional diffusion equation on non-uniform meshes, *J. Comput. Phys.* 265 (5) (2014) 195–210.